

QFT, SOLUTIONS TO PROBLEM SHEET 8

Problem 1: Counterterms in ϕ^4 theory

In ϕ^4 theory in renormalised perturbation theory, show that the Feynman rule for the “2-point vertex” associated to the counterterms δ_Z and δ_{m^2} is

$$\begin{array}{c} p \\ \longrightarrow \\ \text{---}\otimes\text{---} \end{array} = i(p^2\delta_Z - \delta_{m^2}).$$

We have

$$\begin{aligned} \text{---} + \text{---}\otimes\text{---} + \text{---}\otimes\text{---}\otimes\text{---} + \dots &= \text{---} \times \sum_{k=0}^{\infty} \left(\otimes\text{---} \right)^k \\ &= \frac{i}{p^2 + m^2 - i\epsilon} \sum_{k=0}^{\infty} \left(\frac{-p^2\delta_Z + \delta_{m^2}}{p^2 - m^2 + i\epsilon} \right)^k \end{aligned}$$

If this is to converge, one should obtain the usual limit of the geometric series, $\sum_n x^n = (1 - x)^{-1}$, hence

$$\dots = \frac{i}{p^2 - m^2 + i\epsilon} \frac{1}{1 - \frac{-p^2\delta_Z + \delta_{m^2}}{p^2 - m^2 + i\epsilon}} = \frac{i}{p^2 - m^2 + i\epsilon + p^2\delta_Z - \delta_{m^2}} = \frac{i}{Zp^2 - (m^2 + \delta_{m^2}) + i\epsilon}$$

Problem 2: Feynman parameters

Prove the following identities ($A \neq 0$ and $B \neq 0$ are real constants):

1.

$$\frac{1}{AB} = \int_0^1 dx \frac{1}{(xA + (1-x)B)^2}$$

2.

$$\frac{1}{A^n B} = \int_0^1 dx \frac{nx^{n-1}}{(xA + (1-x)B)^{n+1}}$$

3.

$$\frac{1}{A_1 \cdots A_n} = \int_0^1 dx_1 \cdots dx_n \delta(1 - x_1 - \dots - x_n) \frac{(n-1)!}{(x_1 A_1 + \dots + x_n A_n)^n}$$

1. For $A = B$ the identity is obvious. For $A \neq B$ one can directly integrate:

$$\int_0^1 dx \frac{1}{(xA + (1-x)B)^2} = \frac{1}{B-A} \left[\frac{1}{xA + (1-x)B} \right]_{x=0}^{x=1} = \frac{1}{B-A} \left(\frac{1}{A} - \frac{1}{B} \right) = \frac{1}{AB}.$$

2. We differentiate m times on both sides of identity 1:

$$\begin{aligned} (-1)^m m! \frac{1}{A^{m+1}B} &= \frac{d^m}{dA^m} \frac{1}{AB} = \frac{d^m}{dA^m} \int_0^1 dx \frac{1}{(xA + (1-x)B)^2} \\ &= \int_0^1 dx \frac{\partial^m}{\partial A^m} \frac{1}{(xA + (1-x)B)^2} = \int_0^1 dx \frac{(-1)^m x^m (m+1)!}{(xA + (1-x)B)^{2+m}} \end{aligned}$$

The desired relation follows from dividing by $(-1)^m m!$ on both sides and setting $m = n - 1$.

3. This is proved by induction. The statement is true for $n = 2$ (see part 1). Assuming it holds for some given n , let us show that it also holds for $n + 1$:

$$\begin{aligned} \frac{1}{A_1 \dots A_n A_{n+1}} &\stackrel{\text{ass.}}{=} \int_0^1 dy_1 \dots dy_n \delta(1 - y_1 - \dots - y_n) \frac{(n-1)!}{(y_1 A_1 + \dots + y_n A_n)^n} \frac{1}{A_{n+1}} \\ &\stackrel{2.}{=} \int_0^1 dz \int_0^1 dy_1 \dots dy_n \delta(1 - y_1 - \dots - y_n) \frac{n! z^{n-1}}{(zy_1 A_1 + \dots + zy_n A_n + (1-z)A_{n+1})^{n+1}} \end{aligned}$$

The integration domain can be explicitly parametrized as

$$\begin{aligned} &\int_0^1 dz \int_0^1 dy_1 \dots dy_n \delta(1 - y_1 - \dots - y_n) \\ &= \int_0^1 dz \int_0^1 dy_n \int_0^{1-y_n} dy_{n-1} \int_0^{1-y_n-y_{n-1}} dy_{n-2} \dots \int_0^{1-y_n-y_{n-1}-\dots-y_3} dy_2 \Big|_{y_1=1-\sum_{k=2}^n y_k} \end{aligned}$$

Now change variables: call $x_i = zy_i$ ($i = 1 \dots n$) and $x_{n+1} = 1 - z$, so that $dy_i = dx_i/z$ ($i = 2 \dots n$) and $dz = -dx_{n+1}$:

$$\begin{aligned} &\int_0^1 dz \int_0^1 dy_1 \dots dy_n \delta(1 - y_1 - \dots - y_n) \\ &= \int_1^0 (-dx_{n+1}) \int_0^z \frac{dx_n}{z} \int_0^{z-x_n} \frac{dx_{n-1}}{z} \dots \int_0^{z-x_n-x_{n-1}-\dots-x_3} \frac{dx_2}{z} \Big|_{x_1=z-\sum_{k=2}^n x_k} \\ &= \int_0^1 dx_{n+1} \int_0^{1-x_{n+1}} \frac{dx_n}{z} \int_0^{1-x_{n+1}-x_n} \frac{dx_{n-1}}{z} \dots \int_0^{1-x_{n+1}-x_n-\dots-x_3} \frac{dx_2}{z} \Big|_{x_1=z-\sum_{k=2}^n x_k} \\ &= \int_0^1 dx_1 \dots dx_{n+1} \delta(1 - x_1 - \dots - x_{n+1}) z^{1-n} \end{aligned}$$

One finally obtains

$$\frac{1}{A_1 \dots A_n A_{n+1}} = \int_0^1 dx_1 \dots dx_{n+1} \delta(1 - x_1 - \dots - x_{n+1}) z^{1-n} \frac{n! z^{n-1}}{(x_1 A_1 + \dots + x_{n+1} A_{n+1})^{n+1}}$$

as announced.